

Optics

Optics 554 — Homework 1

Fresnel - Fraunhofer diffraction problem.
Due Wednesday, February 10, 2021

Consider Gaussian beam propagation as a diffraction problem.

1. Consider a plane wave with a Gaussian intensity distribution of waist w_0 (as could be obtained from an infinite uniform plane wave, normally incident on an apodized aperture, or simply a Gaussian beam emerging from the flat output coupler of a laser). Write the general Huygens integral for the field at a distance z from this aperture.
2. Consider only one transverse coordinate to find the radial distribution of the electrical field amplitude $\mathcal{E}(x, z)$, using the Fraunhofer approximation.
3. Consider only one transverse coordinate to find the radial distribution of the electrical field amplitude $\mathcal{E}(x, z)$, using the Fresnel approximation.
4. Compare the results of (2) and (3) with that of Gaussian beam propagation, and explain the differences/similarities.

Solutions

Question 1

Each point of the “aperture” is considered to be a source of spherical wave issued from P_0 .

$$E = \frac{E_0 e^{ik\rho(r-r_0)}}{\rho(r-r_0)}. \quad (1)$$

Huygens’s integral is the sum of these spherical waves, with a normalization factor $i/(L\lambda)$, and an “obliquity factor $K[\theta(r, r_0)]$ ”:

$$E(x, y, z) = \frac{i}{L\lambda} \iint_{S_0} \frac{E(x_0, y_0, z_0) e^{-ik\rho(r-r_0)}}{\rho(r-r_0)} K(\theta) dS_0. \quad (2)$$

Question 2

Fraunhofer approximation, in one dimension, is:

$$\begin{aligned} \rho(r-r_0) &\approx \left(z - \frac{k_x}{x}x\right) \\ K(\theta) &= 1 \\ \rho(r-r_0) &= \text{lin the denominator (amplitude term)} \end{aligned} \quad (3)$$

The one-dimensional diffraction pattern is thus given by:

$$E(x, L) \approx \sqrt{\frac{i}{L\lambda}} e^{-ikL} \int_{-\infty}^{\infty} E(x) e^{ik_x x} dx. \quad (4)$$

In terms of the field amplitude; substituting for $E(x, z) \mathcal{E}(x, z) e^{-ikz}$:

$$\mathcal{E}(x, L) = \sqrt{\frac{i}{L\lambda}} \int_{-\infty}^{\infty} \mathcal{E}(x) e^{ik_x x} dx = \sqrt{\frac{i}{L\lambda}} \tilde{\mathcal{E}} k_x. \quad (5)$$

For the initial Gaussian field issued from the beam waist: $\mathcal{E}_0 = \exp(-x^2/w_0^2)$, the Fourier transform is:

$$\tilde{\mathcal{E}}(k_x) = e^{-\frac{k_x^2 w_0^2}{4}} = e^{-\frac{k^2 \theta^2 w_0^2}{4}} = e^{-\frac{k^2 w_0^2 x^2}{4z^2}} = e^{-\frac{x^2}{w^2}}. \quad (6)$$

The transformed beam waist is thus:

$$w^2 = \frac{k^2 w_0^2}{4z^2} = w_0^2 \left(\frac{\lambda}{\pi w_0^2} \right)^2 z^2, \quad (7)$$

which is indeed the expression for a Gaussian beam, in the approximation $z \gg \pi w_0^2/\lambda$.

Question 3

Fresnel approximation, in one dimension, consists in making in the Huygen's integral the substitution:

$$\begin{aligned} K(\theta) &\approx 1 \\ \rho(r - r_0) &\approx (z - z_0) + \frac{(x - x_0)^2 + (y - y_0)^2}{2(z - z_0)}. \end{aligned} \quad (8)$$

The Fresnel integral to solve in one dimension is thus:

$$\begin{aligned} \tilde{\mathcal{E}}(x_2) &= \sqrt{\frac{i}{L\lambda}} \int_{-\infty}^{\infty} \mathcal{E}(x_1) e^{-ik(x_2-x_1)^2/2L} dx_1 \\ &= \sqrt{\frac{i}{L\lambda}} \int_{-\infty}^{\infty} e^{-\frac{x_1^2}{w_0^2}} e^{-\frac{ik(x_2-x_1)^2}{2L}} dx_1. \end{aligned} \quad (9)$$

The result of the integration is also a Gaussian, with a term proportional to x_2^2 in the exponent. Identifying this Gaussian with a real part of $\exp(-x_2^2/w^2)$, we find that w is given by:

$$\frac{1}{w^2} = \text{Re} \left[\frac{ik}{2L} + \frac{k^2}{4L^2} \left(\frac{1}{\frac{1}{w_0^2} + \frac{ik}{2L}} \right) \right]. \quad (10)$$

$$\frac{1}{w^2} = \frac{k^2}{4L^2} \left(\frac{\frac{1}{w_0^2} - \frac{ik}{2L}}{\frac{1}{w_0^2} + \frac{k^2}{4L^2}} \right)$$

This leads to the well known result:

$$w^2 = w_0^2 \left[1 + \left(\frac{\lambda}{\pi w_0^2} \right)^2 L^2 \right] \quad (11)$$