

**Advanced Optics I (463)**  
Test 1, October 15, 2013

## Application of Fresnel and Snell's equations

Consider an interface between air (index 1) and glass (index  $n$ ). A beam of width  $w$  is incident at an angle  $\theta$  on that interface. Demonstrate that the incident power equals the reflected power + the transmitted power (energy conservation). Choose  $s$  or  $p$  polarization (either one should work!).

### Solution

Intensity multiplied by area ( $P = IA$ ) is what is conserved here, so we must account for the change in beam size on transmission. It can be shown using a purely geometric argument (done several times in class) that the relation between the incident  $w_i$  and transmitted  $w_t$  beam waist is,

$$w_t = \frac{\cos \theta_t}{\cos \theta_i} w_i. \quad (1)$$

We also know that intensity is proportional to the electric field squared times the index of refraction:

$$I \sim n |\mathcal{E}|^2. \quad (2)$$

If we set the incident electric field amplitude and incident beam area to 1, then our energy conservation equation becomes:

$$P_i = P_r + P_t = I_r A_r + I_t A_t \quad (3)$$

$$1 = n_i |r|^2 + \frac{n_t \cos \theta_t}{\cos \theta_i} |t|^2, \quad (4)$$

where  $r$  and  $t$  are the fresnel coefficients.

### p-polarization

The Fresnel equations give:

$$\begin{aligned} t_{\parallel} &= \frac{2 \cos \theta}{\cos \theta_t + n \cos \theta} \\ r_{\parallel} &= \frac{\cos \theta_t - n \cos \theta}{\cos \theta_t + n \cos \theta} \end{aligned} \quad (5)$$

Because of the projection of the beam angle on the interface, the energy conservation equation is:

$$R + \frac{n \cos \theta_t}{\cos \theta} T = 1 \quad (6)$$

Substituting:

$$R_{\parallel} + \frac{n \cos \theta_t}{\cos \theta} T_{\parallel} = \frac{\cos^2 \theta_t + n^2 \cos^2 \theta - 2n \cos \theta \cos \theta_t + 4n \cos \theta \cos \theta_t}{(\cos \theta_t + n \cos \theta)^2} = 1 \quad (7)$$

### s-polarization

The Fresnel equations give:

$$\begin{aligned}t_{\perp} &= \frac{2 \cos \theta}{\cos \theta + n \cos \theta_t} \\r_{\perp} &= \frac{\cos \theta - n \cos \theta_t}{\cos \theta + n \cos \theta_t}\end{aligned}\tag{8}$$

Because of the projection of the beam angle on the interface, the energy conservation equation is:

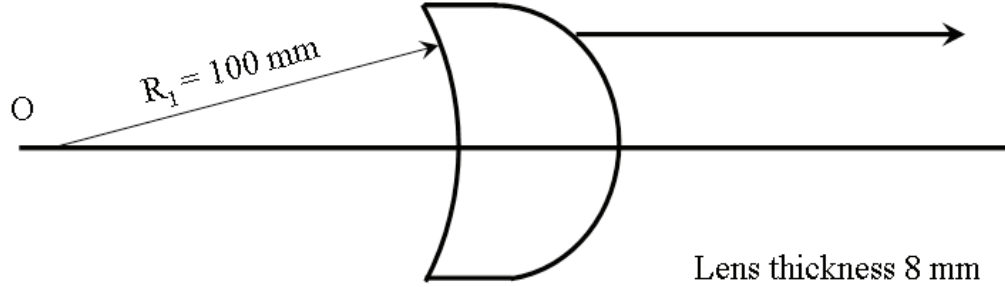
$$R + \frac{n \cos \theta_t}{\cos \theta} T = 1\tag{9}$$

Substituting:

$$R_{\perp} + \frac{n \cos \theta_t}{\cos \theta} T_{\perp} = \frac{\cos^2 \theta + n^2 \cos^2 \theta_t - 2n \cos \theta \cos \theta_t + 4n \cos \theta \cos \theta_t}{(\cos \theta + n \cos \theta_t)^2} = 1\tag{10}$$

## Matrix problem

Several laser cavities are terminated by a curved mirror as output mirror. The rays (inside the laser cavity) are issued from the center of curvature  $O$  (see Fig. 1). It is generally desirable that the output beam be collimated. Given a radius of curvature of the first surface  $R = 100$  mm, a thickness  $t = 8$  mm, calculate within the paraxial approximation, the curvature of the second surface in order to have a collimated output. The index of refraction is  $n = 1.5$ . Hint: try to solve with the least amount of matrices possible (you do not need to start from point  $O$ ).



What is the curvature of the second surface?

Figure 1:

## Solution

In the paraxial approximation ( $\sin \alpha \approx \alpha$ ), the angle that each ray makes with the horizontal axis is  $\alpha = y/R_1$ . Since any ray issued from  $O$  is orthogonal to the spherical input surface of the lens, there is no refraction at the curved surface, and we can use as initial matrix inside the glass:

$$\begin{pmatrix} y \\ \frac{y}{R_1} \end{pmatrix}.$$

All that is left to do is to translate by the thickness  $t$ , and the second curved interface:

$$\begin{pmatrix} y' \\ \alpha' \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -\frac{n-1}{R_2} & n \end{pmatrix} \begin{pmatrix} 1 & t \\ 0 & 1 \end{pmatrix} \begin{pmatrix} y \\ \frac{y}{R_1} \end{pmatrix} = \begin{pmatrix} 1 & t \\ -\frac{n-1}{R_2} & n - \frac{n-1}{R_2}t \end{pmatrix} \begin{pmatrix} y \\ \frac{y}{R_1} \end{pmatrix}$$

The definition of a collimated beam is that  $\alpha' = 0$  which leads to:

$$0 = -(R_1 + t)\frac{n-1}{R_2} + n$$

leading to the result:

$$R_2 = \frac{n-1}{n}(t + R_1) = \frac{0.5}{1.5} \times 108 = 36$$

## Prism problem

A beam is incident from the left on the prism sketched in Fig. 2. The index of refraction of the prism is 2.1. It is immersed in a medium of which you vary the index from 1 to infinity.

- Complete the figure indicating the path of the beam inside the prism, and the value of the angles of incidence on the various faces.
- Assuming the beam is polarized with the electric in the plane of the figure, calculate characteristic values of the index of refraction
- Using the characteristic values calculated in (b), sketching (qualitatively) a graph with the values of the intensity transmitted through the face  $AB$  versus the index of refraction (*qualitatively*: i.e. it is not necessary to use Fresnel formulae)
- What is the value of the intensity transmitted through the face  $OB$ , when the outside index takes the values of 1.05 and for 2.1?

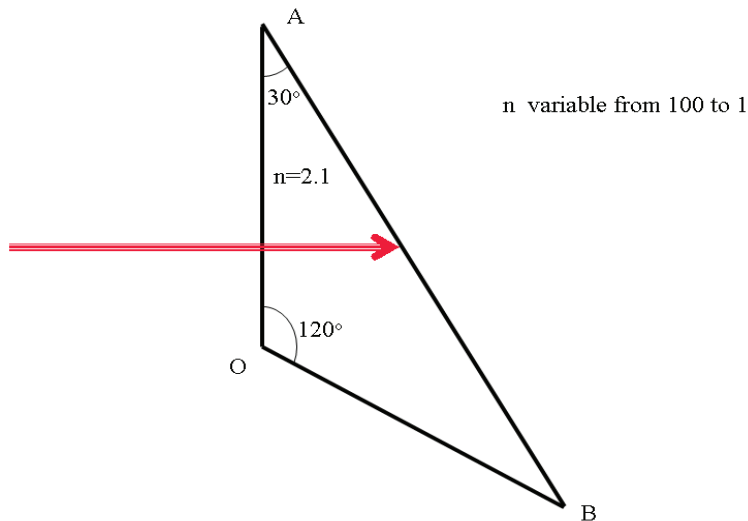


Figure 2:

## Solution

**a**

Because the initial angle of incidence is  $0^\circ$ , we can conclude that the angle between the input beam and  $\overline{AB}$  is  $90 - 30 = 60^\circ$ . Which means the angle of incidence at the second interface ( $\overline{AB}$ ) must be  $30^\circ$ . Using the law of reflection ( $\theta_i = \theta_r$ ), we know that the reflected beam will then also have an angle of incidence with the third interface  $\overline{OB}$  of  $0^\circ$ .

**b**

A beam polarized in the plane of the figure has TM, P, or  $\parallel$  polarization. This means the characteristic values will occur at TIR ( $n=1.05$ ), Brewster's angle ( $n=1.21$ ), and when the index of refraction on the outside of the prism is the same as the inside ( $n=2.1$ ).

**c**

Below the TIR condition ( $n=1.05$ ), there will be no power transmitted through interface  $\overline{AB}$ . After  $n=1.05$ , the transmitted power will increase to a local maximum at the Brewster's angle of  $n=1.21$ . Note that the transmission will not be 1 here due to the fresnel loss at the first interface. As  $n$  keeps increasing the transmitted power will drop slightly before increasing back up to the maximum transmission of 1 at  $n=2.1$ . All power is transmitted here since at  $n=2.1$  there is essentially no prism. As  $n$  increases to infinity, the transmitted power will gradually drop due to the fresnel reflections at the interfaces.

**d**

At  $n=1.05$ , the ray will be totally internally reflected at  $\overline{AB}$ , and therefore the maximum amount of power will be transmitted through  $\overline{OB}$ . The Intensity will not be 1, however, due to the fresnel reflections at  $\overline{AO}$  and  $\overline{OB}$ . The loss at these interfaces will be,

$$r_{\parallel} = \frac{n_t - n_i}{n_t + n_i} = \frac{2.1 - 1.05}{2.1 + 1.05} = 0.333. \quad (11)$$

Which means the power through interface  $\overline{OB}$  when  $n=1.05$  will be 0.778.

When  $n=2.1$ , the beam will encounter no interfaces and therefore 0 intensity will be transmitted through  $\overline{OB}$ .